

# Gauss' law in presence of dielectrics

We write,  $\mathbf{D} \equiv \epsilon_0 \mathbf{E} + \mathbf{P}$

└─ Electric Displacement

→ In terms of  $\mathbf{D}$ , Gauss's law reads

$$\nabla \cdot \mathbf{D} = \rho_f$$

Hence, the integral form:

$$\oint \mathbf{D} \cdot d\mathbf{a} = Q_{f_{\text{enc}}}$$

→  $Q_{f_{\text{enc}}}$  denotes the total free charge enclosed in the volume.

*This is particularly useful since we only need to know the free charge in the system, which is something we can control. The bound charges appear only after polarization. In a typical problem, initially we just start with the free charge densities.*

# Linear dielectrics

*Linear dielectrics is a class of dielectrics for which the induced polarization is proportional to the applied electric field.*

$$\mathbf{P} = \epsilon_0 \chi_e \mathbf{E}.$$

$\chi_e$   $\longrightarrow$  **Electric susceptibility** of the medium; depends on microscopic structure of the substance in question and also on external conditions such as temperature.

In linear media we have

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} = \epsilon_0 \mathbf{E} + \epsilon_0 \chi_e \mathbf{E} = \epsilon_0 (1 + \chi_e) \mathbf{E},$$

so  $\mathbf{D}$  is *also* proportional to  $\mathbf{E}$ :

$$\mathbf{D} = \epsilon \mathbf{E},$$

where

**(Dielectric Constant)**  $\epsilon_r \equiv 1 + \chi_e = \frac{\epsilon}{\epsilon_0}$

$$\epsilon \equiv \epsilon_0 (1 + \chi_e).$$

Permittivity of free space  $\epsilon_0$   
Permittivity of a material  $\epsilon$

# Linear dielectrics: Example

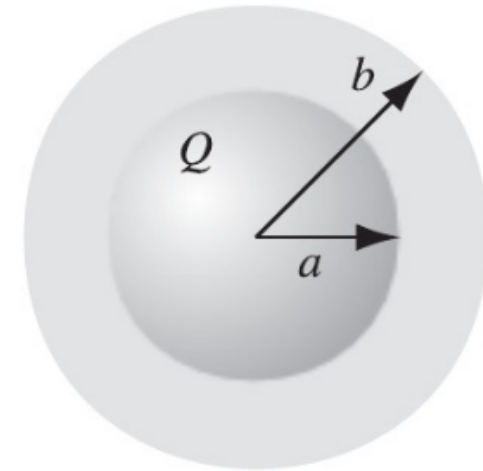
A metal sphere of radius  $a$  carries a charge  $Q$ . It is surrounded, out to radius  $b$ , by linear dielectric material of permittivity  $\epsilon$ . Find the potential at the center (relative to infinity).

→ let's begin by calculating  $\mathbf{D}$ ,

$$\mathbf{D} = \frac{Q}{4\pi r^2} \hat{\mathbf{r}}, \quad \text{for all points } r > a.$$

Then,

$$\mathbf{E} = \begin{cases} \frac{Q}{4\pi \epsilon r^2} \hat{\mathbf{r}}, & \text{for } a < r < b, \\ \frac{Q}{4\pi \epsilon_0 r^2} \hat{\mathbf{r}}, & \text{for } r > b. \end{cases}$$



# Linear dielectrics: Example

The potential at the center is therefore

$$\begin{aligned} V &= - \int_{\infty}^0 \mathbf{E} \cdot d\mathbf{l} = - \int_{\infty}^b \left( \frac{Q}{4\pi\epsilon_0 r^2} \right) dr - \int_b^a \left( \frac{Q}{4\pi\epsilon r^2} \right) dr - \int_a^0 (0) dr \\ &= \frac{Q}{4\pi} \left( \frac{1}{\epsilon_0 b} + \frac{1}{\epsilon a} - \frac{1}{\epsilon b} \right). \end{aligned}$$

→ To find the bound charges:

$$\mathbf{P} = \epsilon_0 \chi_e \mathbf{E} = \frac{\epsilon_0 \chi_e Q}{4\pi \epsilon r^2} \hat{\mathbf{r}},$$

in the dielectric, and hence

$$\rho_b = -\nabla \cdot \mathbf{P} = 0,$$

while

$$\sigma_b = \mathbf{P} \cdot \hat{\mathbf{n}} = \begin{cases} \frac{\epsilon_0 \chi_e Q}{4\pi \epsilon b^2}, & \text{at the outer surface,} \\ -\frac{\epsilon_0 \chi_e Q}{4\pi \epsilon a^2}, & \text{at the inner surface.} \end{cases}$$

(unit vector points outward with respect to the dielectrics)

